



# Effect of Velocity in Icing Scaling Tests

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## EFFECT OF VELOCITY IN ICING SCALING TESTS

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### Abstract

This paper presents additional results of a study first published in 1999 to determine the effect of scale velocity on scaled icing test results. Reference tests were made with a 53.3-cm-chord NACA 0012 airfoil model in the NASA Glenn Icing Research Tunnel at an airspeed of 67 m/s, an *MVD* of 40  $\mu\text{m}$ , and an *LWC* of 0.6 g/m<sup>3</sup>. Temperature was varied to provide nominal freezing fractions of 0.8, 0.6 and 0.5. Scale tests used both 35.6- and 27.7-cm-chord 0012 models for <sup>2</sup>/<sub>3</sub>- and <sup>1</sup>/<sub>2</sub>-size scaling. Scale test conditions were found using the modified Ruff (AEDC) scaling method with the scale velocity determined in five ways. Four of the scale velocities were found by matching the scale and reference values of water-film thickness, velocity, Weber number and Reynolds number. The fifth scale velocity was simply the average of those found by matching the Weber and Reynolds numbers. The resulting scale velocities ranged from 85 to 220% of the reference velocity. For a freezing fraction of 0.8, the value of the scale velocity had no effect on how well the scale ice shape simulated the reference shape. For nominal freezing fractions of 0.5 and 0.6, the best simulation of the reference shape was achieved when the scale velocity was the average of the constant-Weber-number and the constant-Reynolds-number velocities.

### Nomenclature

$A_c$	Accumulation parameter, dimensionless
$b$	Relative heat factor, dimensionless
$c$	Model chord, m
$c_{p,ws}$	Specific heat of water at the surface temperature, cal/g K
$h_{film}$	Water-film thickness at leading edge, m
$K_0$	Modified inertia parameter, dimensionless
<i>LWC</i>	Cloud liquid-water content, g/m <sup>3</sup>
$M$	Mach number, dimensionless
<i>MVD</i>	Water droplet median volume diameter, $\mu\text{m}$

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$n$	Freezing fraction, dimensionless
$r_{le}$	Leading-edge radius of airfoil, m
$Re$	Reynolds number of model, dimensionless
$Re_\delta$	Reynolds number of water droplet, dimensionless
$t_{st}$	Static temperature, °C
$t_{tot}$	Total temperature, °C
$V$	Air velocity, m/s
$We$	Weber number based on droplet size and water properties, dimensionless
$\phi$	Droplet energy transfer terms in energy balance, °C
$\Lambda_f$	Latent heat of freezing, cal/g
$\rho_i$	Ice density, g/m <sup>3</sup>
$\theta$	Air energy transfer terms in energy balance, °C
$\tau$	Accretion time, min

### Introduction

This paper presents the results of a continuation of a 1998 study published in 1999<sup>1</sup> to determine the effect of scale velocity on scaled ice shapes.

To minimize test-section blockage in icing tunnels, it is often necessary to test reduced-size models. Scaling methods are thus required to determine scaled test conditions that will simulate the full-scale icing encounter and produce the same non-dimensional ice shape. Similarity parameters have been identified which best describe the important phenomena of icing physics. Scaling methods<sup>2,3</sup> have been developed by selecting those similarity parameters that appear to have the most influence on ice accretion, and requiring that the parameters have the same value for both scale and reference tests. The phenomena important to icing include the flowfield approaching and around the model, the droplet trajectories, the quantity of ice accumulation, the surface heat balance and, possibly, phenomena related to the dynamics of surface water on the model. For rime ice, because water freezes on impact, only the first three of these phenomena affect the ice shape.

The flowfield can be simulated by matching scale and reference values of  $Re$  and  $M$ , by using a model whose

non-dimensional coordinates are the same as the full-scale (reference) article and by setting the scale angle of attack equal to the reference. For icing encounters, the speeds involved are usually low such that  $M$  should have little effect and is neglected. In the past,  $Re$  has usually been ignored by arguing that any ice accretion will trip the boundary layer and the flow will then be independent of  $Re$ . However, Bilanin<sup>4</sup> advocated including  $Re$  in icing scaling analyses, and recently, the study of reference 1 used  $Re$  as an optional scaling parameter.

Similarity of droplet trajectories and, therefore, droplet collection efficiencies, can be obtained by matching the modified inertia parameter,  $K_0$ , of Langmuir and Blodgett.<sup>5</sup>

Similarity of ice accumulation results from a match of the accumulation parameter,  $A_c$ :

$$A_c = \frac{LWC V \tau}{2 r_{le} \rho_i} \quad (1)$$

Here,  $r_{le}$  is the leading-edge radius of the model. The length scale used to compute the values of all of the similarity parameters in this study was twice the leading-edge radius. For the NACA 0012, the leading-edge radius is 1.58% of the chord.

The energy balance at the surface is based on the work of Messinger.<sup>6</sup> The energy balance can be written in the form,

$$n = \frac{c_{p,ws}}{\Lambda_f} \left( \phi + \frac{\theta}{b} \right) \quad (2)$$

where  $n$  is the freezing fraction;  $\phi$ , the water energy transfer parameter;  $\theta$ , the air energy transfer parameter; and  $b$ , the relative heat factor, defined by Tribus, et al.<sup>7</sup> Equation (2) provides additional similarity parameters for scaling, but only three of the four are independent. For this study, only  $n$ ,  $\phi$  and  $\theta$  will be considered.

Finally, some understanding of surface dynamics can provide still more similarity parameters. Two of these are the Weber number,  $We$ , and the non-dimensional water-film thickness,  $h_{film}/c$ . Kind<sup>8</sup> has also identified parameters, but for the present study, these will not be considered. For this investigation,  $We$  was based on the droplet size. An experimental correlation for water-film thickness was given by Feo and Urdiales:<sup>9</sup>

$$\frac{h_{film}}{c} \propto LWC^{1/2} (Re)^{-1/4} (Re_\delta)^{15/4} (We)^{-9/4} \quad (3)$$

This correlation was based on tests of the effect of rain on aircraft performance; thus, the water droplets were at least an order of magnitude larger than those experienced in conditions described by FAA FAR 25 Appendix C. The  $LWC$ 's in the Feo and Urdiales study were also higher than those of interest here. Consequently, equation (3) may not be applicable to Appendix C icing. Because there is no other data relating water film thickness to icing parameters, however, equation (3) will be used with reservations.

From the above discussion, then, eight scaling parameters can be considered:  $Re$ ,  $K_0$ ,  $A_c$ ,  $n$ ,  $\phi$ ,  $\theta$ ,  $We$  and  $h_{film}/c$ . After the reference conditions for the test to be simulated have been selected along with the scale-to-reference size ratio, there are five test conditions for the scale test that need to be established: the static temperature, velocity, cloud droplet size, cloud liquid-water content and the icing time. In icing tunnels with altitude capability, the scale test-section pressure is a sixth condition that must be determined. The present work was performed in a sea-level tunnel, so pressure was established by ambient conditions.

Ruff<sup>3</sup> performed tests in the AEDC R-1D icing tunnel to evaluate the importance of the parameters,  $K_0$ ,  $A_c$ ,  $n$ ,  $\phi$ , and  $\theta$  to scaling. He found that the scaled ice shapes agreed best with reference shapes when all five parameters were matched to the reference values. This approach to scaling is known as the Ruff (or AEDC) Method. Because the R-1D allows altitude simulation, Ruff's work included the calculation of the scale test-section pressure in addition to temperature,  $MVD$ ,  $LWC$  and time. The scale velocity is selected arbitrarily in the Ruff Method. The Ruff Method has also been used in sea-level tunnels in a modified form in which  $\theta$  is ignored and just the remaining four parameters are matched. Scale velocity is chosen by the user for this modified Ruff Method, as well.

In Ruff's experiments, the scale velocity was often simply equated with the reference value. This is a practical approach, and it coincidentally insures that  $\theta$ ,  $b$  and  $M$  will match between the scale and reference situation. Although some more recent studies have been done with velocity selected by matching  $We$ <sup>10,11</sup>, until the study of reference (1), the choice of scale velocity and its effect on scaling results has not received much attention.

Reference (1) reported a series of 1/2-size scaling tests performed in the NASA Glenn Icing Research Tunnel. Five sets of reference conditions were tested with the full-size model (53.3-cm-chord NACA 0012 airfoil). Tests with a 26.7-cm-chord 0012 were made with scale conditions determined by the modified Ruff method.

(Because the IRT is a sea-level tunnel, the parameter  $\theta$  was ignored.) Five possible scale velocities were determined for each reference case. These were found by matching either  $h_{film}/c$ , velocity,  $We$ ,  $Re$ , or the average of velocities found by matching  $We$  and  $Re$ . This average- $V$  method has no physical basis, but it was used to assist in spotting possible trends with velocity. The matching of  $h_{film}/c$  was referred to as the Feo method; it produced a scale velocity of about 85% of the reference. The scale velocity for the constant- $We$  approach was about 130%, for the average- $V$ , about 175%, and for the constant- $Re$ , about 220% of the reference. The scale velocities for the constant- $Re$  method in some cases were higher than the tunnel capabilities permitted and, in others, produced total temperatures above freezing. Even the average- $V$  method sometimes gave total temperatures above freezing. Thus, not all of the five sets of reference cases could be tested with all scale velocities.

The study of reference (1) found that for tests with freezing fractions higher than 0.8, including some tests with rime ice, scale velocity had little effect on scaled shapes. For a freezing fraction of about 0.5, there was better agreement between the scale and reference ice shapes as the scale velocity increased, as long as the scale total temperature was below about  $-2^{\circ}\text{C}$ . At a freezing fraction of 0.3 a clear trend could not be established since even the average-velocity method gave a total temperature above  $0^{\circ}\text{C}$ .

Additional tests were performed in 1999 to obtain supplementary data and to investigate some of the questions raised by the 1998 study reported in reference (1). Because that previous study ignored the scaling parameter  $\theta$ , the present work included some tests to evaluate whether results would be different if  $\theta$  were used instead of  $\phi$  as part of the system of equations to determine scale conditions. While constant  $\phi$  results in simpler equations to solve, constant  $\theta$  produces a lower scale total temperature. In addition to the  $1/2$ -scale results, a few sets of ice shapes will be shown for scaling from a chord of 53.3 cm to 35.6 cm ( $2/3$  scale). NACA 0012 models were again tested in the NASA Glenn Icing Research Tunnel at  $0^{\circ}$  angle of attack. Along with results from the 1999 tests, this paper will also present additional results from 1998.

#### Reference Conditions

All reference tests were made with 53.3-cm-chord models. The nominal test conditions are shown in table I. Three reference cases, designated 110, 111 and 112, were taken from the study of reference (1). Each case has a different freezing fraction due to the different

Table I. Reference Test Conditions  
Reference Model, 53.3-cm-Chord NACA 0012  
Reference Airspeed, 67 m/s  
Reference  $MVD$ ,  $40\text{ }\mu\text{m}$   
Reference  $LWC$ ,  $0.6\text{ g/m}^3$   
Reference Spray Time, 11.2 min

Case	$t_{sts}$ $^{\circ}\text{C}$	$t_{tot}$ $^{\circ}\text{C}$	$n$
110	-13.9	-11.7	0.8
111	-11.1	-8.9	0.7
112	-8.4	-6.1	0.5

temperature. The airspeed, droplet size, liquid-water content and spray time were the same for all cases. The reference shapes from reference (1) were used again for this study. Actual recorded test conditions for both reference and scale tests will be given with the ice shapes. The values of some of the scaling parameters to be given in this paper may differ from those previously reported due to changes in the calculation routines used.

#### Test Description

##### NASA Glenn Icing Research Tunnel

The facility and hardware used are shown in figure 1. Figure 1(a) is a plan view of the NASA Glenn Icing Research Tunnel (IRT).<sup>12,13</sup> The IRT has a test section width of 2.7 m and a height of 1.8 m. It has a refrigeration system that allows accurate control of the test-section temperature from  $-40$  to  $4^{\circ}\text{C}$ . A water spray system with ten spray bars simulates the conditions in a natural icing cloud. The test-section cloud droplet size,  $MVD$ , and liquid-water content,  $LWC$ , depend on spray-bar air and water pressures. The relationships among these pressures, the tunnel airspeed and the cloud properties are established periodically by a series of tunnel calibration tests.<sup>14</sup> The cloud has been calibrated over a range of test-section airspeeds from 22 to 156 m/s and droplet median volume diameters from 14 to  $50\text{ }\mu\text{m}$ . Two sets of spray nozzles, the Mod-1 and Standard, are used to provide different ranges of liquid-water content. Depending on the nozzle set, the airspeed and the droplet diameter, the test-section liquid-water content can be controlled from less than  $0.2$  to over  $5\text{ g/m}^3$ . The Standard nozzles were used for the 1999 tests, while the Mod-1's were used in 1998.

The IRT spray system includes a water valve at each nozzle, and water is recirculated prior to the start of the spray. Spray-bar air and water pressures can thus be

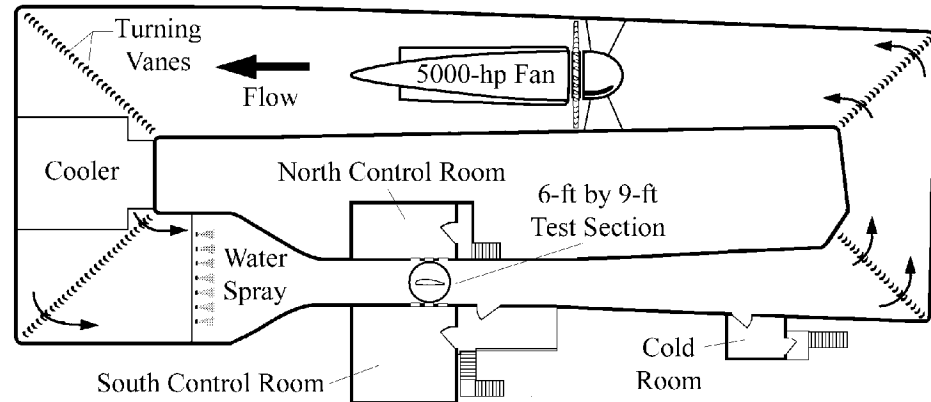
established and stabilized before spray is admitted to the tunnel. This feature virtually eliminates start-up transients that were present previously in the IRT.

A cold room is mounted on the side of the tunnel with access through the diffuser as well as from the outside. The cold-room temperature can be controlled to permit ice studies off line. A three-dimensional scanner in the cold room is available for digitally recording the ice-shape coordinates of ice accretions. Although the ice produced in this study was scanned, the shapes to be presented here were obtained from two-dimensional hand tracings of the ice cross section.

#### Test Models and Procedures

The same NACA 0012 airfoil models were used in both the 1998 and 1999 studies. Chords were 53.3 cm (reference model), 35.6 cm and 26.7 cm. Each was mounted vertically, at  $0^\circ$  angle of attack, in the center of the IRT test section. The models were machined from solid aluminum. All had a span of 61 cm and were placed between end plates as shown in figure 1(b). The models' mid spans were located on the tunnel centerline (midway between floor and ceiling of the test section). The five horizontal lines marked around the leading edge of the model indicate the mid-span and positions 2.5 and 5 cm above and below mid-span. Ice shapes were recorded only at the mid-span and 5-cm-above locations for the 1999 study and at mid-span,  $\pm 2.5$  cm and 5 cm above mid-span for the 1998 tests. In this paper, only the mid-span shapes will be shown.

The test temperature and airspeed were first established, then the spray-bar air and water pressures were set and stabilized. To initiate the spray, the water valves were opened and the spray timer started. At the completion of the desired time, the spray was turned off and the tunnel fan stopped. For the 1999 tests, the model was removed from its stand and carried to the cold room where the surface coordinates were recorded with the 3-D laser scanner. After scanning, a thin heated aluminum plate was used to melt horizontal slots through the accreted ice at specific span-wise locations. The ice shapes were traced by hand onto cardboard templates inserted into these slots in the ice. Two



(a) NASA Glenn Icing Research Tunnel (IRT) Flow Loop and Features.



(b) 53.3-cm-Chord NACA 0012 Model in IRT Test Section.

Figure 1. Test Facility and Model.

models of each chord were fabricated so that while the ice accretion was being scanned and traced in the cold room on the first model, a second model could be installed in the test section and tested. Finally, the model in the cold room was cleaned. When the test with the second model was completed, the clean model was reinstalled in the test section and the procedure repeated. The 1998 procedures were described in reference (1). They were similar to those described above except that most ice shapes were recorded from the model in the test section.



Following the test, the ice-shape tracings were digitized to provide a permanent record of each shape. The ice shapes to be shown in this report were plotted from these 2-D digital records.

#### Average Test Conditions

Tunnel and cloud conditions were recorded at 1-sec intervals for tests with the 26.7- and 35.6-cm-chord models and at 2-sec intervals for the 53.3-cm-chord model. Averages of the 1- or 2-sec readings over the spray period were made for each instrument used. The test-section total temperature at any moment was the average of the readings of eleven type-T thermocouples located on the turning vanes upstream of the spray bars. The test-section velocity was found from the averages of the pressures from two pitot-static probes on opposite walls at the entrance to the test section. The *MVD* and *LWC* were calculated from the averages of the ten spray-bar air pressures and water-air differential pressures. The conditions reported in this paper are these instrument averages averaged over time. These conditions will be tabulated as part of the ice-shape comparisons in figures 3 to 6. The scaling parameters shown in these figures have been computed from these average test conditions.

#### Uncertainty

Estimates of the uncertainty in the reported average conditions were made by considering fluctuations of the values with time, possible instrument errors including calibration, uncertainties in tunnel calibration of *MVD* and *LWC*, and differences in measurements from one location to another in the test section. From this analysis for the conditions of the tests reported here total temperatures are probably known to within  $\pm 1.5^\circ\text{C}$ , velocity,  $\pm 2.5\%$ , *MVD*,  $\pm 11\%$  and *LWC*,  $\pm 12\%$ . There is no absolute standard for drop-sizing instruments; therefore, the use of different instruments or operators to measure the clouds tested could produce *MVD*'s outside this range of uncertainty.

The uncertainties in the scaling parameters were determined from the above test-condition uncertainties.  $K_0$  is thus estimated to be known within about  $\pm 22\%$ ,  $A_c \pm 12\%$ ,  $n \pm 10\%$ ,  $b \pm 12\%$ ,  $\phi \pm 3\%$ ,  $\theta \pm 3\%$ ,  $Re \pm 3\%$ ,  $We \pm 12\%$  and  $M \pm 3\%$ . At this time, the effect on the ice shape of varying scaling parameters by these magnitudes is unknown.

### Results and Discussion

#### Considerations of Nozzle Type and Span-wise Position

Some of the ice-shape comparisons to be presented will be made between tests made in 1998 and 1999 using

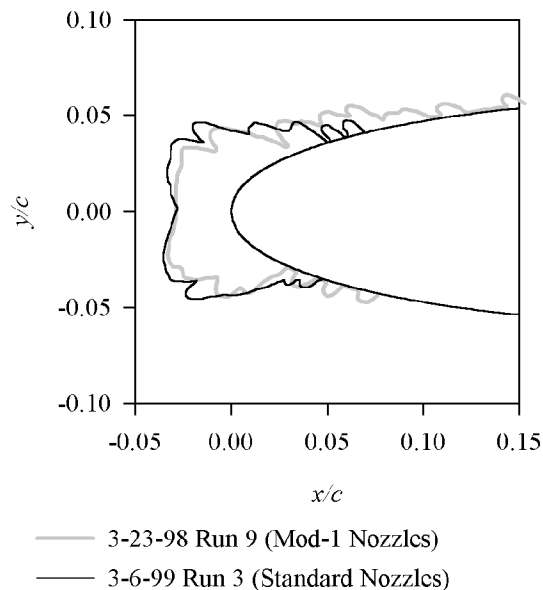


Figure 2. Comparison of Results From Tests Using Mod-1 and Standard Nozzles.  $c$ , 26.7 cm;  $t_{st}$ ,  $-11^\circ\text{C}$ ;  $V$ , 88 m/s; *MVD*, 24  $\mu\text{m}$ ; *LWC*, 0.77  $\text{g/m}^3$ ;  $\tau$ , 3.7 min.

different nozzle sets. To determine consistency between the results for the two types of nozzles, some tests from 1998 were repeated in 1999. Figure 2 compares the results of a scale test from 1998 using the Mod-1 nozzle set with one run in 1999 at the same conditions but with the Standard nozzles. The two ice accretions displayed only minor differences, and these were typical of the shape variations observed when icing sprays are repeated with the same nozzle. The quality of the shape agreement indicated by figure 2 was observed with several other repeat tests that were made. It was concluded that the 1998 and 1999 results can be directly compared.

Reference (1) showed comparisons of ice shapes traced at different span-wise locations from 2.5 cm below the mid span to 5 cm above. No significant variations in ice shape were seen. Similarly, in the tests made in 1999, no meaningful differences between the mid-span shapes and those recorded 5 cm above mid span were observed. Thus, small errors in locating the span-wise positions for the tracings should have had no effect on ice shape.

#### 1/2-Size Scaling with Rime Ice

When there is sufficient convective heat transfer to absorb all the latent heat released, water freezes immediately on impact to form rime ice. Consequently, a description of the heat balance at the surface is unnecessary, and, with no liquid water on the ice,

dynamics of a liquid surface are irrelevant. For this situation, only the reference accumulation and inertia parameters,  $A_c$  and  $K_0$ , have to be matched by the scale test, and any method to choose scale velocity should be equally valid. Reference (1) presented results of scaling to  $\frac{1}{2}$  size with scale velocities using the Feo, constant-velocity, constant- $We$  and average-velocity methods. The four scale tests produced evidence that rime shapes should not change with scale velocity. However, some rime ice shapes in reference (1) were smaller than expected due to an inconsistency in the tunnel  $LWC$  calibration over the range of velocities tested. That calibration has since been revised, and the 1999 tests used the latest IRT calibration. No additional rime results will be shown here.

#### $\frac{1}{2}$ -Size Scaling at Nominal Freezing Fraction of 0.8 (fig. 3)

The ice shapes are shown in the plots with non-dimensional coordinates. The reference shape is drawn with a gray line and is repeated for comparison with the scaled shapes in each portion of the figure. Figure 3(a) shows the scaled shape which resulted from testing with the Feo scale velocity; figure 3(b), the constant-velocity; figure 3(c), the constant- $We$ ; and figure 3(d), the average of the constant- $We$  and constant- $Re$  velocities. The ice shapes shown in figures 3(a), (b) and (c) have already been reported in reference (1). The average-velocity result in figure 3(d) adds a higher velocity scale test than what was given in reference (1). The first column of numbers below the plots of ice shapes in figure 3(a) gives the reference test conditions and scaling parameters. The scale test conditions are listed under the corresponding portion of the figure.

There are usually three attributes of an ice shape that need to be evaluated when judging the effectiveness of scaling. They are the quantity of ice accreted, the limits of accretion, and the size and location of significant features such as horns. The scale velocity had little effect on any of these characteristics of the scale ice shape at this freezing fraction, although the Feo velocity, just 85% of the reference, gave a flatter leading region than the shapes formed at higher velocities. None of the shapes exactly reproduced the leading-edge features. The accumulation parameter,  $A_c$ , varied by about  $\pm 7\%$  from test to test, and the quantity of ice appeared to be insensitive to this variation.

#### $\frac{1}{2}$ -Size Scaling at Nominal Freezing Fraction of 0.6 (fig. 4)

Figure 4 gives scaling results for constant velocity (fig. 4(a)), constant  $We$  (fig. 4(b)), average velocity (fig. 4(c)), and constant  $Re$  (fig. 4(d)). As with results

discussed above, all were made with the water energy transfer parameter,  $\phi$ , held constant. A fifth scale test was made with the  $Re$  again matched, but with the air energy transfer parameter,  $\theta$ , instead of  $\phi$ , matched to the reference value (fig. 4(e)). The use of  $\theta$  in place of  $\phi$  results in a scale total temperature that is significantly reduced. Reference (1) noted that one of the practical problems of using constant- $Re$  scaling is that the total temperature can be near or above freezing. Temperatures near freezing were found to produce scaled ice shapes that had little resemblance to reference shapes.

The relative sizes of the scaled shapes appear to be slightly larger than the reference, but that may have been because the scale accumulation parameter was as much as 9% higher than desired. Accretion limits are sometimes difficult to define precisely from traced ice shapes, but for the scale results of figure 4, there appeared to be little effect of scale velocity on the accretion limit of the main ice shape. Scale velocity had a noticeable effect on the horn position, however.

Increasing scale velocity from 67 m/s (constant velocity, fig. 4(a)) to 87 m/s (constant  $We$ , fig. 4(b)) produced only minor differences in the scale horn position. When the velocity was increased to 117 m/s (average velocity, fig. 4(c)), the main scale horn angle and size conformed closely to the reference. The scale feathers also matched in size, although in the reference case, feathers were observed further aft than in the scale test.

The two constant- $Re$  scaled tests (Figs. 4(d) and (e)), produced main ice shapes which approximately matched the reference size and shape in the horn region, but failed to simulate the size or extent of the feathers which formed aft of the main shape in the reference test. Reference (1) speculated that the high scale velocity which results when  $Re$  is matched may strip small feathers from the surface early in their formation, thus preventing substantial feather growth from taking place. It is not evident from the results of figure 4 that there is any advantage in using scale velocities higher than the average of those obtained by applying constant  $We$  and constant  $Re$ .

There does not seem to be a significant difference between the scale shapes formed using either constant  $\phi$  (fig. 4(d)) or constant  $\theta$  (fig. 4(e)) when  $Re$  is matched. This suggests that either of these parameters may be equally effective.

#### $\frac{1}{2}$ -Size Scaling at Nominal Freezing Fraction of 0.5 (fig. 5)

The nominal freezing fraction for case 112 was 0.5. The reference ice shape featured the horns typical of glaze

ice with small feathers farther back on the airfoil. Figure 5 shows the five scaled ice shapes that resulted from testing with scale velocities ranging from those for the Feo method to the constant- $Re$  method. It can be seen from figure 5 that all scale velocities produced approximately the correct quantity of ice, and accretion limits appeared to be nearly the same as for the reference shape. However, as the scale velocity increased, the horns on the scaled shapes moved forward. The best match of the scaled and reference horn shape was obtained when the scale velocity was the average of the constant- $We$  and constant- $Re$  velocities. The horns for the constant- $Re$  scaling projected forward of the reference horns.

The 1999 scale test results for the Feo, constant- $V$ , constant- $We$ , and average- $V$  methods given in figure 5 were in close agreement with the 1998 results reported in reference 1. For the 1999 tests, the constant- $Re$  conditions were found by matching  $\theta$  since the constant- $\phi$  approach, used in 1998, produced a scale total temperature slightly above freezing. The use of constant  $\theta$  produced a scaled shape (fig. 5(e)) which, although not a perfect match, was more consistent with the reference shape than was the constant- $\phi$  scaled shape of 1998.<sup>1</sup>

The table of conditions below the plotted shapes indicates that the scaled tests had accumulation parameters as much as 9% higher than the reference and freezing fractions as much as 17% lower than the reference. These variations are greater than desired, so caution should be observed when forming conclusions.

#### $\frac{2}{3}$ -Size Scaling at Nominal Freezing Fraction of 0.5 (fig. 6)

Scaling size to only  $\frac{2}{3}$  the reference permits testing with constant  $Re$  at more moderate scale velocities than are needed for  $\frac{1}{2}$ -scale testing. A series of tests were made scaling from 53.3-cm to 35.6-cm chord, using the same reference case as for the tests shown in figure 5 (case 112, with a nominal freezing fraction of 0.5). Five scale velocities were again tested, using the Feo, constant-velocity, constant- $We$ , average velocity and constant- $Re$  methods. The constant- $Re$  scaled conditions were determined with  $\phi$  matched to the reference value.

The results are shown in figure 6. Because the scale size was closer to the reference, the differences between the scale and reference shapes are not as pronounced as those seen in figure 5. The relative quantity of ice and accretion limits for the scaled tests matched the reference for all scale velocities. However, the same effect of scale velocity on ice shape can be seen in both

figures 5 and 6, with the scale horns moving forward as velocity increased. The best match of scale and reference ice shapes was again found for the average velocity method, and the scale horns were again somewhat too far forward when constant  $Re$  was applied.

For the scale tests of figure 6, the accumulation parameter varied from the reference by less than 2% and the freezing fraction by about 4%. Because of the improved parameter matching over that in figure 5, the figure 6 shape comparisons may have more credibility. In any case, the conclusions from the two figures are consistent: (1) The scale velocity has little effect on quantity of ice or accretion limits, and (2) the best match of scaled horn size and position was obtained when the scale velocity was chosen as the average of the constant- $We$  and the constant- $Re$  velocities.

#### Concluding Remarks

Scaling tests were performed in the NASA Glenn Icing Research Tunnel using a 53.3-cm-chord NACA 0012 model for reference tests with 35.6- and 26.7-cm-chord scale models. Scale test conditions were found using the Ruff scaling method for sea-level tunnels along with five methods to select scale velocity. Four of these methods were based on insuring that some physically based parameter had the same value for both scale and reference conditions. The first method matched the Feo water film thickness, the second, the test-section velocity, the third, the Weber number, the fourth, the Reynolds number. The fifth method averaged the velocities found by matching the Weber number and by matching the Reynolds number. These methods gave scale velocities that ranged from about 85% of the reference to about 220% for  $\frac{1}{2}$ -scale testing and from 90% to 155% for the  $\frac{2}{3}$ -scale tests. For some tests the effect of matching the air-energy transport parameter,  $\theta$ , instead of the water-energy transport parameter,  $\phi$ , was evaluated.

The observations and conclusions from this study are:

1. Limited testing showed no effect of whether constant  $\phi$  or constant  $\theta$  was used in the scaling equations. Constant  $\phi$  uses simpler calculation routines, but constant  $\theta$  gives a lower scale total temperature when the scale velocity is higher than the reference.
2. The quantity of ice was adequately simulated by any of the scale velocities tested. Accretion limits were difficult to define precisely, but appeared to be approximately simulated using all of the scale velocities considered.

3. For a freezing fraction of about 0.8, there appeared to be no effect of scale velocity on scale ice shape. This observation is consistent with the results of reference (1).
4. For nominal freezing fractions of 0.6 and 0.5, the scale horns moved forward as scale velocity increased. The best scaling resulted from averaging the constant- $We$  and the constant- $Re$  scale velocities. For a freezing fraction of 0.5, the constant- $Re$  method produced ice shapes whose horns projected forward beyond those of the reference shape.

One-half-scale tests with nominal freezing fractions of 0.6 and 0.5 and two-thirds-scale tests with a nominal freezing fraction of 0.5 all indicated that average-velocity scaling gave the best simulation of the reference ice shape. This method of choosing the scale velocity has no physical basis, but its success could possibly indicate that both  $We$  and  $Re$  play nearly equally important roles in the physics of ice accretion. Thus, while these two parameters cannot be simultaneously matched using a practical scaling method, the average-velocity approach may provide an acceptable compromise. It is also possible that some other physical parameter is coincidentally satisfied, at least approximately, by the constant-velocity method. In any case, before final conclusions can be reached, tests with additional reference conditions and different scaling ratios are needed to confirm the observations reported here.

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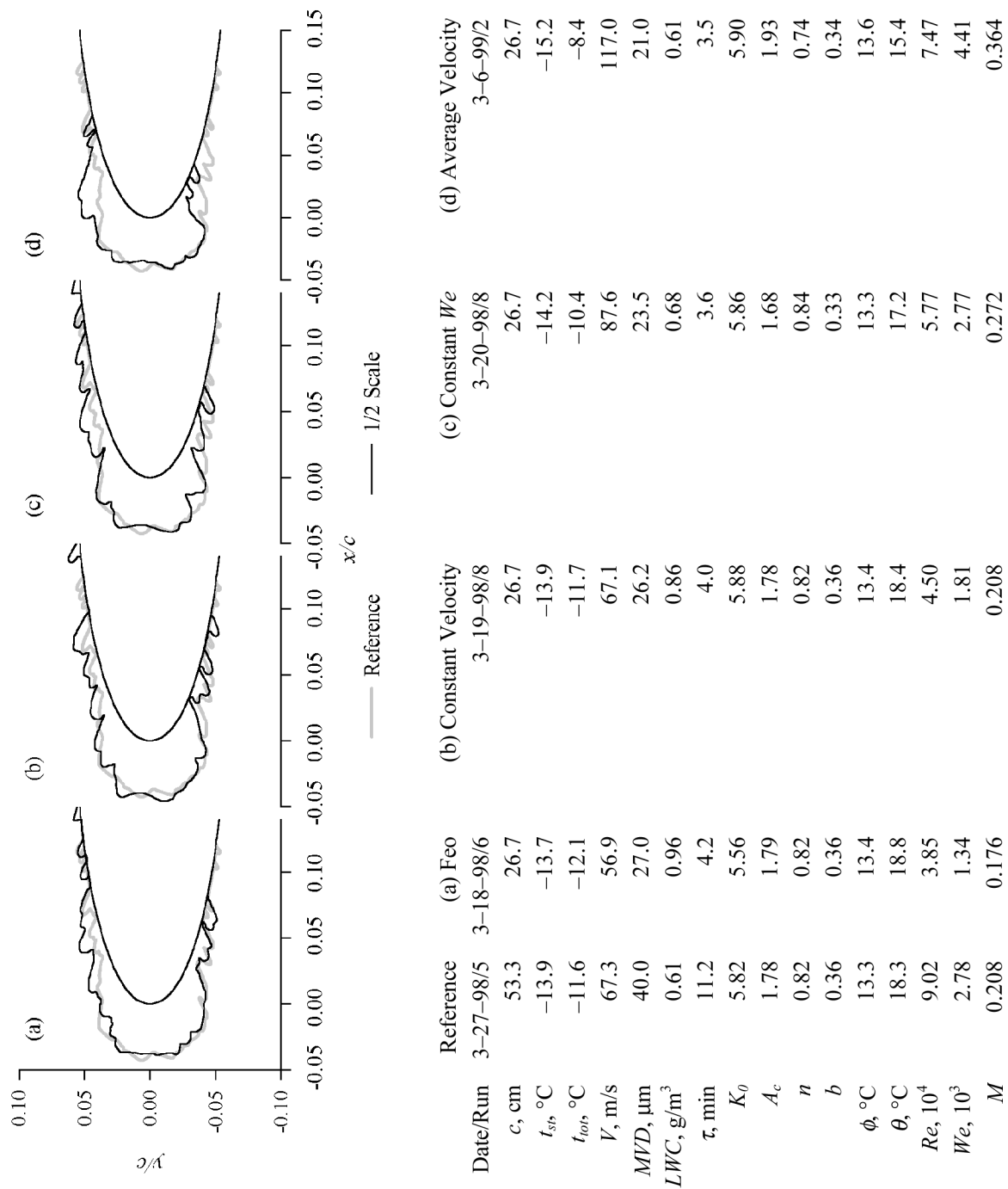
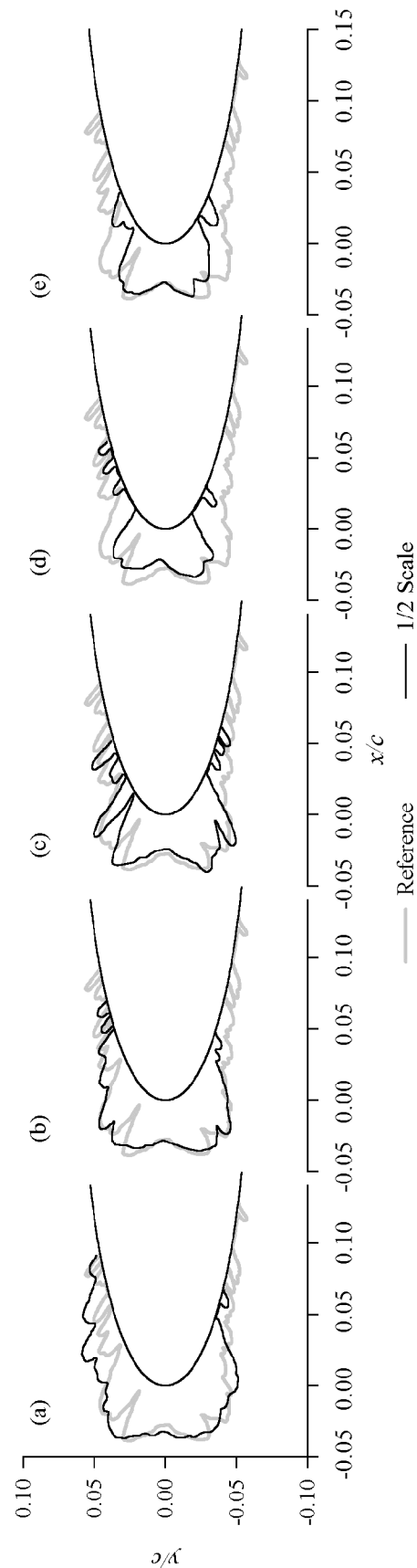
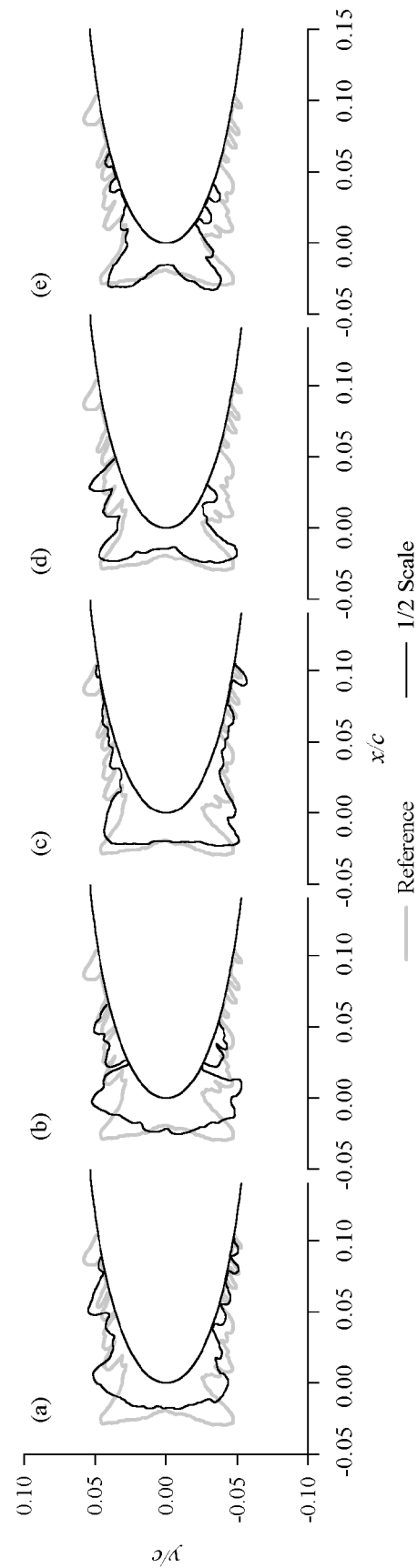


Figure 3. Half-Size Scaling for Case 110. All Scale Results Compared to Same Reference Shape.



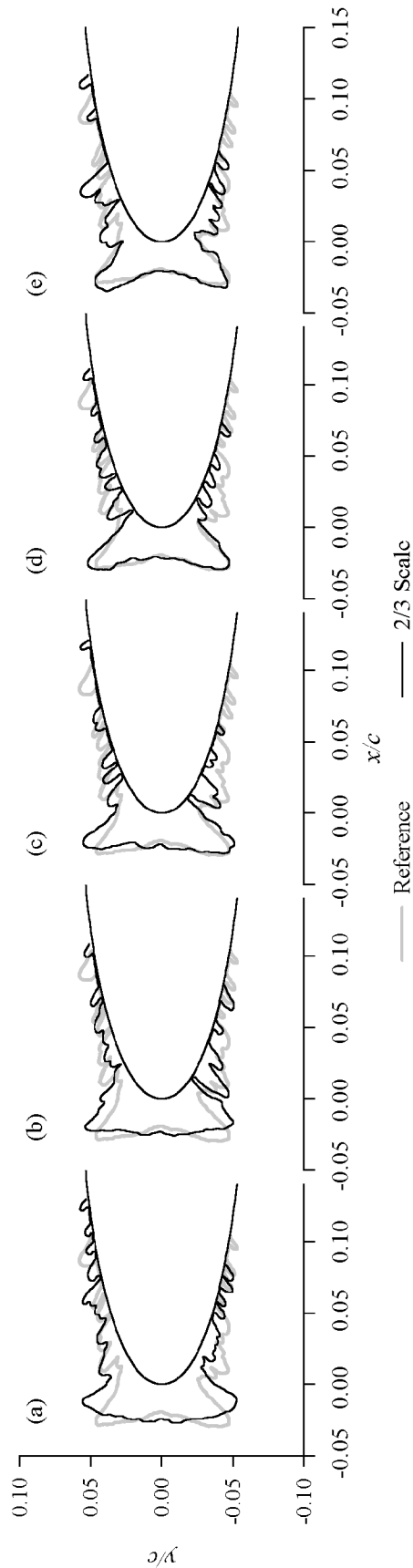
	Reference	(a) Const $V$	(b) Constant $We$	(c) Average Velocity	(d) Constant $Re$	(e) Const. $Re$ , Const. $\theta$
Date/Run	3-26-98/6	3-4-99/4	3-6-99/3	3-6-99/9	3-23-98/4	3-6-99/4
$c$ , cm	53.3	26.7	26.7	26.7	26.7	26.7
$t_{st}$ , °C	-11.1	-10.9	-11.3	-12.0	-13.2	-17.0
$t_{tot}$ , °C	-8.9	-8.7	-7.5	-5.1	-2.3	-7.0
$V$ , m/s	67.1	67.2	87.6	117.0	147.4	141.8
$MVD$ , $\mu\text{m}$	40.0	26.3	23.7	20.9	19.0	19.1
$LWC$ , g/m <sup>3</sup>	0.61	0.93	0.77	0.59	0.39	0.68
$\tau$ , min	11.2	4.0	3.7	3.6	3.8	2.5
$K_0$	5.80	5.92	5.93	5.89	5.90	5.79
$A_c$	1.77	1.92	1.93	1.92	1.69	1.88
$n$	0.65	0.60	0.59	0.56	0.57	0.61
$b$	0.36	0.39	0.37	0.33	0.25	0.43
$\phi$ , °C	10.6	10.4	10.4	10.3	10.6	14.6
$\theta$ , °C	14.7	14.4	13.3	11.2	8.8	14.6
$Re$ , 10 <sup>4</sup>	8.84	4.42	5.66	7.31	8.85	8.81
$We$ , 10 <sup>3</sup>	2.77	1.83	2.79	4.41	6.35	5.89
$M$	0.207	0.207	0.270	0.361	0.456	0.442

Figure 4. Half-Size Scaling for Case 111. All Scale Results Compared to Same Reference Shape.



	Reference	(a) Feo	(b) Constant Velocity	(c) Constant $We$	(d) Average Velocity	(e) Const. $Re$ , Const. $\theta$
Date/Run	3-26-98/4	3-6-99/5	3-6-99/8	3-5-99/10	3-5-99/4	3-5-99/8
$c$ , cm	53.3	26.7	26.7	26.7	26.7	26.7
$t_{stb}$ , °C	-8.4	-8.1	-8.1	-8.5	-9.4	-13.8
$t_{tot}$ , °C	-6.1	-6.5	-5.9	-4.8	-2.6	-3.7
$V$ , m/s	67.0	56.3	66.7	87.2	117.3	142.2
$MVD$ , $\mu\text{m}$	40.0	28.0	26.3	23.6	21.0	19.4
$LWC$ , g/m <sup>3</sup>	0.61	1.04	0.93	0.75	0.56	0.72
$\tau$ , min	11.2	4.2	4.0	3.7	3.8	2.4
$K_0$	5.80	5.85	5.90	5.87	5.92	5.94
$A_c$	1.77	1.91	1.93	1.88	1.93	1.90
$n$	0.48	0.45	0.44	0.43	0.40	0.43
$b$	0.36	0.40	0.39	0.36	0.32	0.46
$\phi$ , °C	7.8	7.7	7.6	7.6	7.8	11.4
$\theta$ , °C	10.8	11.1	10.5	9.5	7.6	10.4
$Re$ , 10 <sup>4</sup>	8.67	3.67	4.31	5.54	7.21	8.65
$We$ , 10 <sup>3</sup>	2.76	1.37	1.80	2.76	4.44	6.02
$M$	0.205	0.173	0.204	0.268	0.360	0.441

Figure 5. Half-Size Scaling for Case 112. All Scale Results Compared to Same Reference Shape.



Date/Run	Reference	(a) Feo	(b) Constant Velocity	(c) Constant $We$	(d) Average Velocity	(e) Constant $Re$
	3-26-98/4	3-25-98/8	3-25-98/7	3-24-98/5	3-24-98/4	3-24-98/3
$c$ , cm	53.3	35.6	35.6	35.6	35.6	35.6
$t_{ss}$ , °C	-8.4	-8.2	-8.2	-8.4	-8.8	-9.1
$t_{tot}$ , °C	-6.1	-6.3	-6.0	-5.3	-4.6	-3.7
$V$ , m/s	67.0	61.0	67.1	78.2	91.7	104.3
$MVD$ , $\mu\text{m}$	40.0	32.5	31.5	29.5	27.5	26.5
$LWC$ , g/m <sup>3</sup>	0.61	0.80	0.75	0.66	0.57	0.50
$\tau$ , min	11.2	6.3	6.1	5.9	5.7	5.7
$K_0$	5.80	5.87	5.93	5.88	5.82	5.96
$A_c$	1.77	1.78	1.78	1.77	1.74	1.74
$n$	0.48	0.48	0.47	0.46	0.47	0.46
$b$	0.36	0.36	0.36	0.35	0.32	0.31
$\phi$ , °C	7.8	7.7	7.7	7.7	7.8	7.8
$\theta$ , °C	10.8	11.0	10.6	10.1	9.5	8.7
$Re$ , 10 <sup>4</sup>	8.67	5.28	5.78	6.68	7.74	8.69
$We$ , 10 <sup>3</sup>	2.76	1.86	2.18	2.77	3.55	4.43
$M$	0.205	0.187	0.206	0.240	0.281	0.320

Figure 6. Two-Thirds-Size Scaling for Case 112. All Scale Results Compared to Same Reference Shape.



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